

Charge and the electric field

About 200 years ago, people noticed that *charged* objects exerted a force on each other.



The force was observed to be:

1. Directed along a line between the two objects
2. Proportional to the product of the charges Qq
3. Proportional to the *inverse square* of the distance between them

We can write the force as a vector. In the coordinate system centred on the positive charge, we have

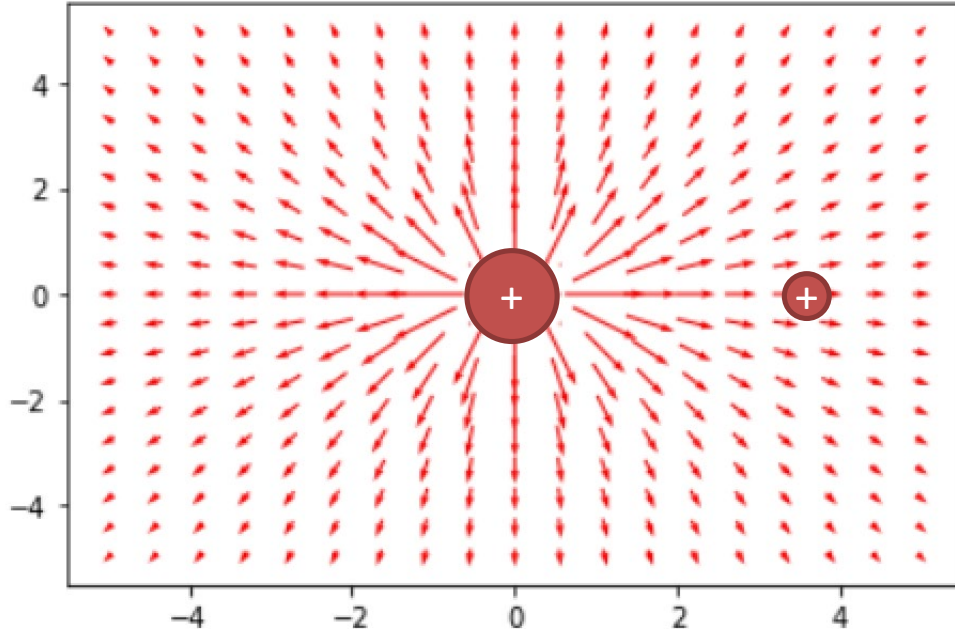
$$\mathbf{F} = \frac{Qq}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\mathbf{r}}$$

ϵ_0 is an experimentally-measured number, called the vacuum *permittivity*:

$$\epsilon_0 = 8.854188 \times 10^{-12} \text{ A}^2 \text{ s}^2 / \text{N} / \text{m}^2$$

The *electric charge* is measured in *Coulombs* C.

The current understanding* is that small electric *charges* emit a *vector field*. This field in turn exerts a force on other charges.



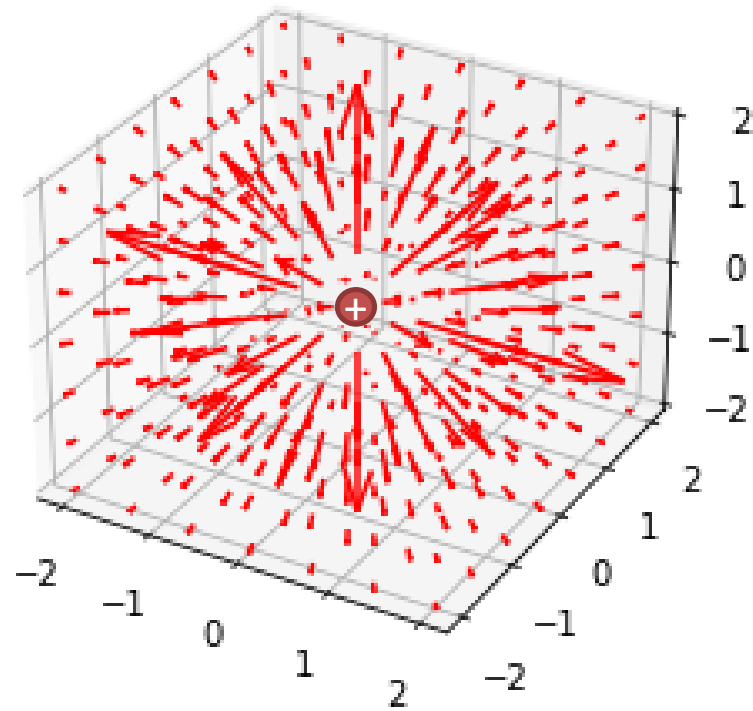
$$\mathbf{F} = \frac{Qq}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\mathbf{r}}$$
$$= q\mathbf{E}$$

The electric field from a charge Q is

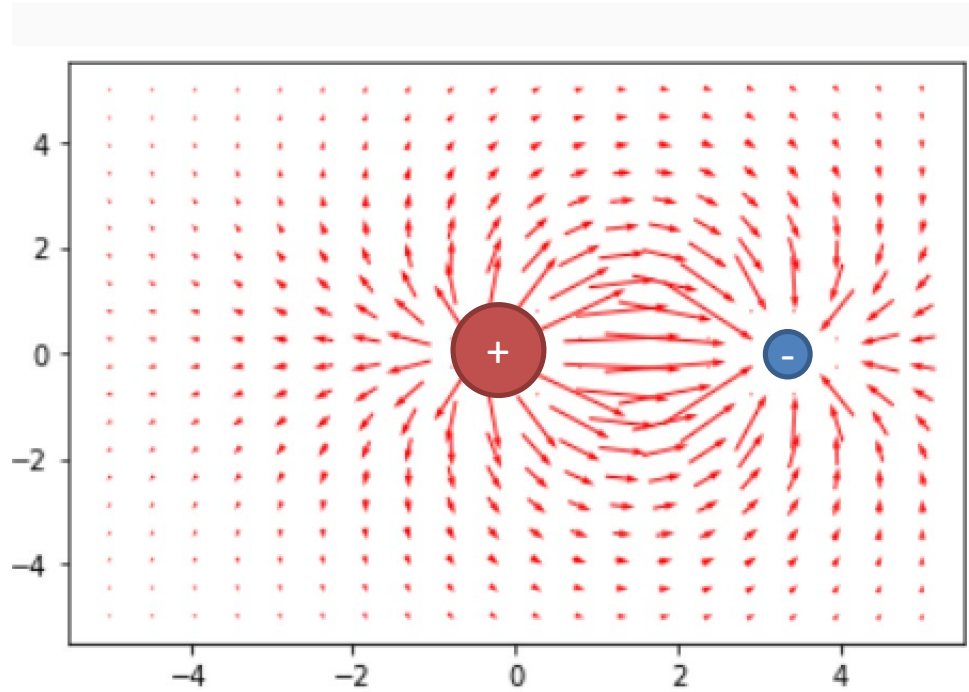
$$\mathbf{E} = \frac{Q}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\mathbf{r}}$$

*of classical theory; quantum theory has its own fields!

In 3D the electric field can look a bit more complicated:



Meanwhile, each of the *other charges* emits its own contribution to the field.

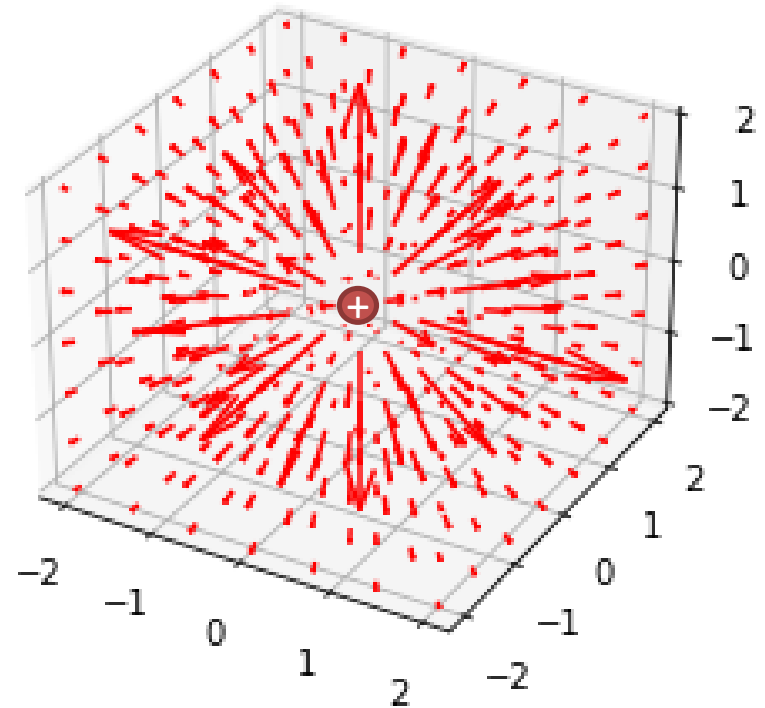


$$\mathbf{E} = \frac{Q_1}{4\pi\epsilon_0} \frac{1}{|\mathbf{r}_1|^2} \hat{\mathbf{r}}_1 + \frac{Q_2}{4\pi\epsilon_0} \frac{1}{|\mathbf{r}_2|^2} \hat{\mathbf{r}}_2$$

We consider a single “point charge” charge q , centred at the origin. The $\mathbf{E}$ field is

$$\mathbf{E} = \frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\mathbf{r}}$$

We now compute the *flux integral* of $\mathbf{E}$ over a spherical surface of radius R :

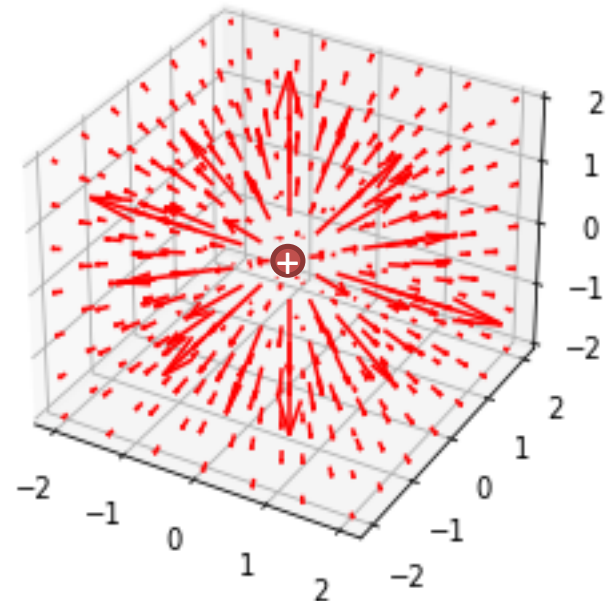


The resulting integral does not depend on the radius, or (it turns out) on the shape of the surface. In general we have

$$\iint_S \mathbf{E} \cdot d\mathbf{S} = \frac{Q_{\text{enc}}}{\epsilon_0}$$

where S is *any surface enclosing the charge*.

What happens when we apply the divergence theorem?



So we have found

$$\iint_S \mathbf{E} \cdot d\mathbf{S} = \iiint_V (\nabla \cdot \mathbf{E}) dV = \frac{Q_{\text{enc}}}{\epsilon_0}$$

But we can write the total enclosed charge in terms of the *charge density* $\rho(\mathbf{r})$:

The only way this works for any volume V is if

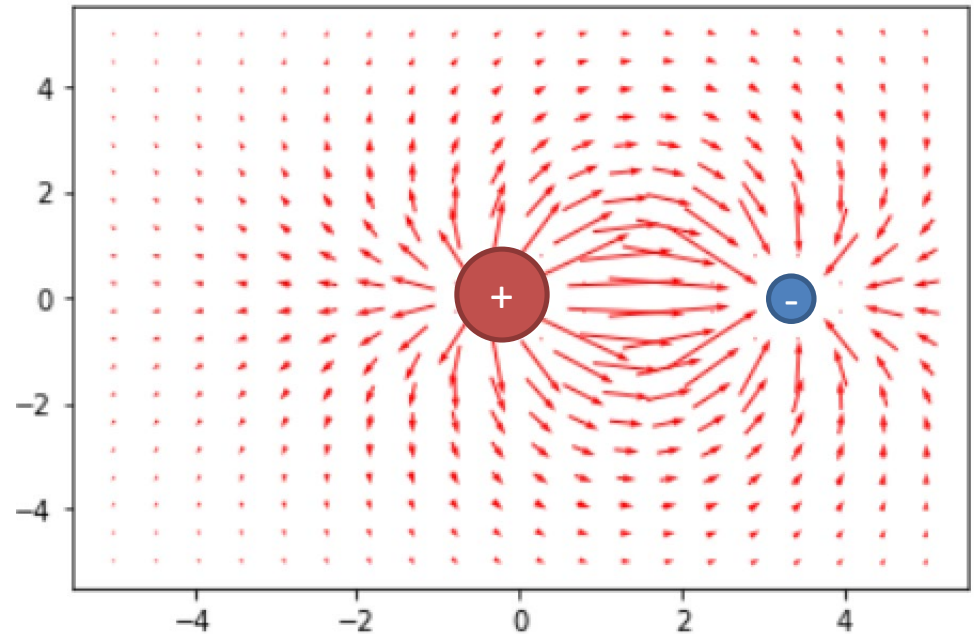
$$\nabla \cdot \mathbf{E} = \frac{\rho(\mathbf{r})}{\epsilon_0}$$

This is known as *Coulomb's law* and it is the first of Maxwell's equations.

What does Coulomb's law tell us?

$$\nabla \cdot \mathbf{E} = \frac{\rho(\mathbf{r})}{\epsilon_0}$$

1. Positive charge is a *source* of electric field, while negative charge is a *sink*.



2. The flux integral

$$\iint_S \mathbf{E} \cdot d\mathbf{S} = \frac{Q_{\text{enc}}}{\epsilon_0}$$

only depends on the enclosed charge – charges outside the surface do not contribute anything.

Example: Compute the electric field for a sphere of radius R , having a uniform charge density ρ .

Example: A continuous line of charge can be represented by a *line charge density (charge per unit length)* $\lambda(\mathbf{r})$. Compute the Electric field for an infinite line charge aligned along the z axis.

What about the curl of the electric field?

Consider the field from a point charge:

$$\mathbf{E} = \frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\mathbf{r}}$$

So the Electric field is irrotational:

$$\nabla \times \mathbf{E} = \mathbf{0}$$

We can therefore express it as the gradient of a potential function:

By convention

$$\mathbf{E} = -\nabla V$$

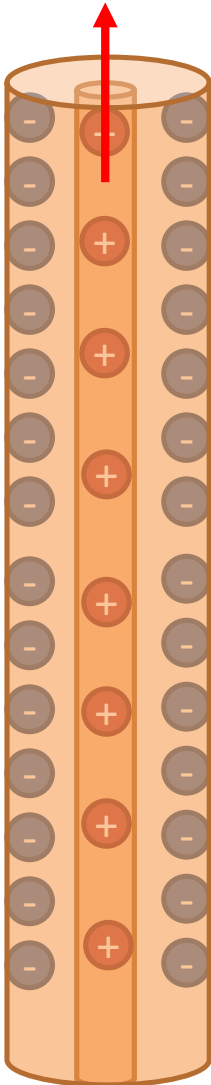
where V is known as the electrostatic potential. Taking the divergence leads to

$$\nabla \cdot (-\nabla V) = \frac{\rho}{\epsilon_0}$$

Electric current and the magnetic field

An electric current is a *moving line of charges*, in which *positive and negative charges are moving continuously in opposite directions*

Positive
charge
flow



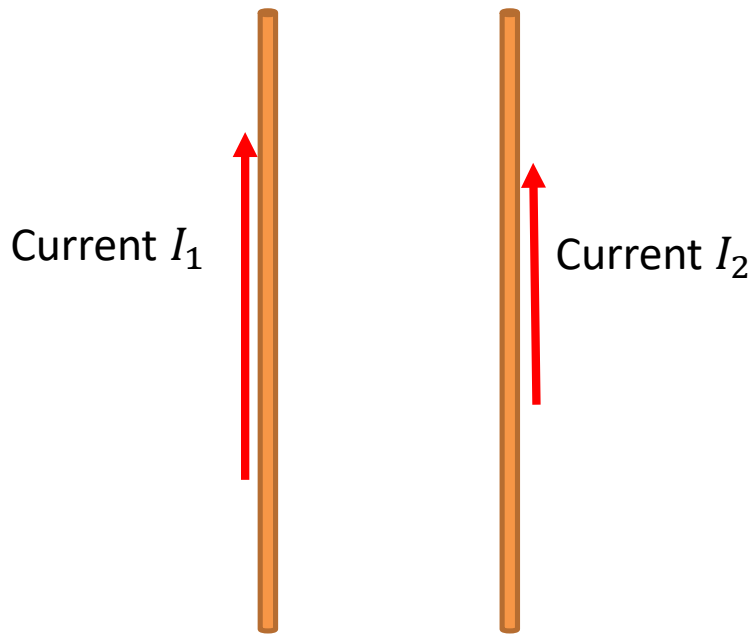
Negative
charge
flow

We can depict the current
as a vector I , which gives
the net flow of positive charge per second

Current I_1



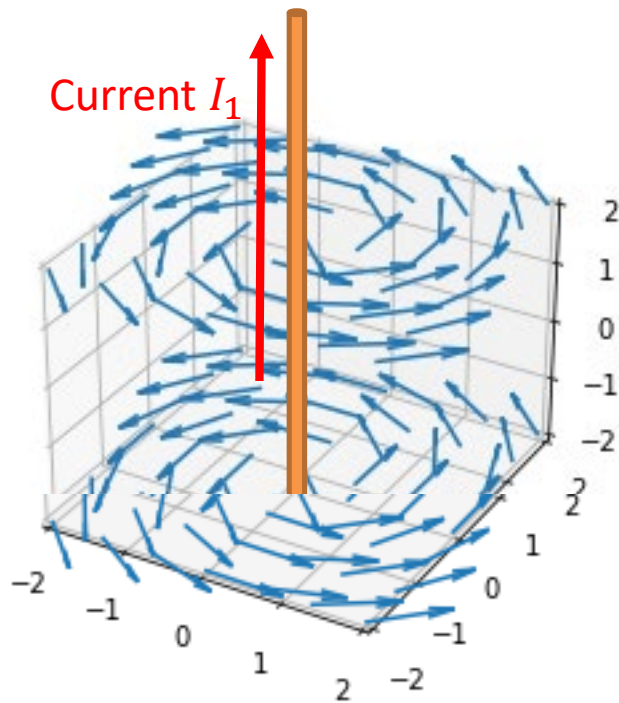
Observation: electric currents exert forces on each other



The force per unit length is

1. Directed along a line between the two currents
2. Proportional to the product of the currents $I_1 I_2$
3. *Inversely proportional* to the distance between them

Each current emits a *magnetic field B*, which circles around the *z*-axis:



$$\mathbf{B} = \frac{\mu_0}{2\pi} \frac{I_1}{r} \hat{\boldsymbol{\theta}}$$

And the resulting force per unit length on *another current* is

$$\mathbf{f} = \mathbf{I}_2 \times \mathbf{B}$$

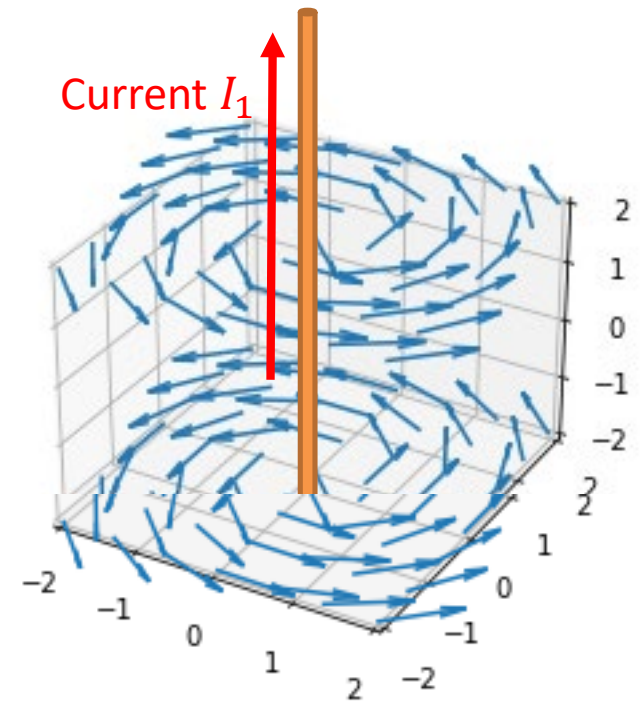
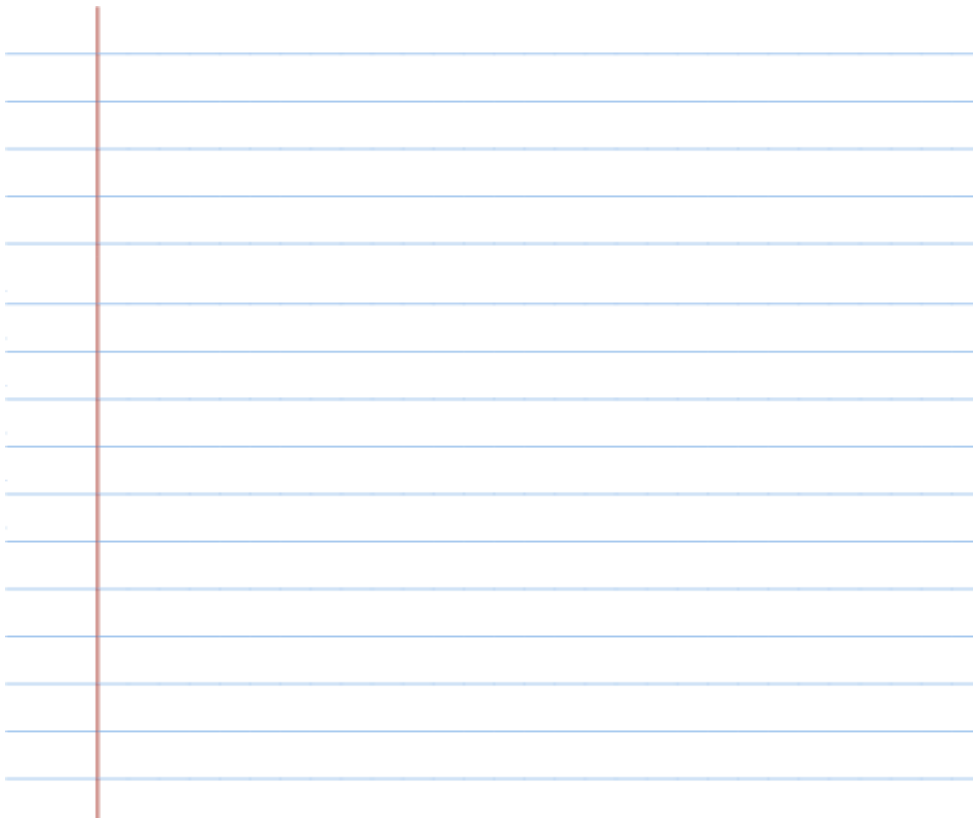
The constant μ_0 can be measured experimentally, and is called the *vacuum permeability*

$$\mu_0 = 1.256637 \times 10^{-6} \text{N/A}^2$$

We now take the *vector line integral* of

$$\mathbf{B} = \frac{\mu_0}{2\pi} \frac{I_1}{r} \hat{\boldsymbol{\theta}}$$

along a loop of radius R enclosing the wire:



So

$$\oint \mathbf{B} \cdot d\mathbf{r} = \mu_0 I_{\text{enc}}$$

Where I_{enc} is the enclosed charge. Note that this does not depend on the radius of the loop.

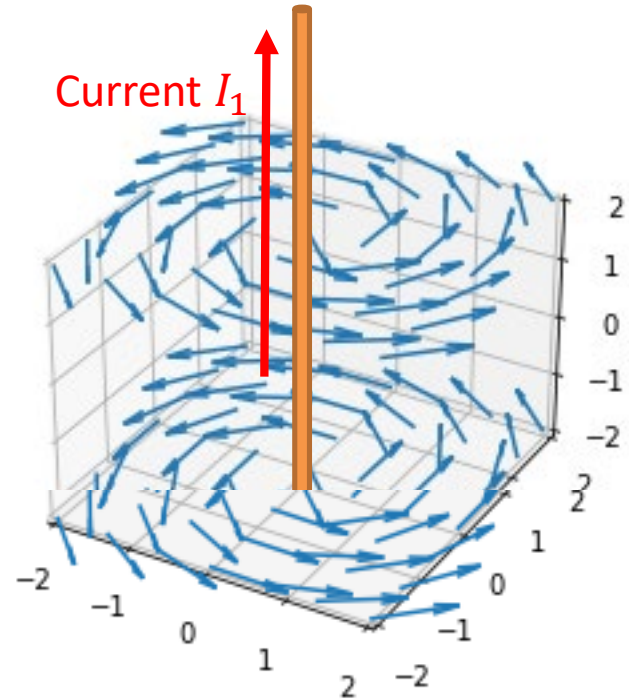
By Stokes' theorem

$$\oint \mathbf{B} \cdot d\mathbf{r} = \iint_S \nabla \times \mathbf{B} \cdot d\mathbf{S} = \mu_0 I_{\text{enc}}$$

But we can also write the current in terms of *the current density* $\mathbf{J}(\mathbf{r})$:

$$I_{\text{enc}} = \iint_S \mathbf{J} \cdot d\mathbf{S}$$

Therefore:



We have found an expression for the curl of the magnetic field:

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$$

What about the divergence?

From the expression for a line current

$$\mathbf{B} = \frac{\mu_0}{2\pi} \frac{I_1}{r} \hat{\boldsymbol{\theta}}$$

The divergence is

We therefore have the following equations for the magnetic field:

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$$

Maxwell's equations

What have we found so far?

For the electric field

$$\nabla \cdot \mathbf{E} = \frac{\rho(\mathbf{r})}{\varepsilon_0}$$

$$\nabla \times \mathbf{E} = \mathbf{0}$$

For the magnetic field

$$\nabla \cdot \mathbf{B} = 0$$

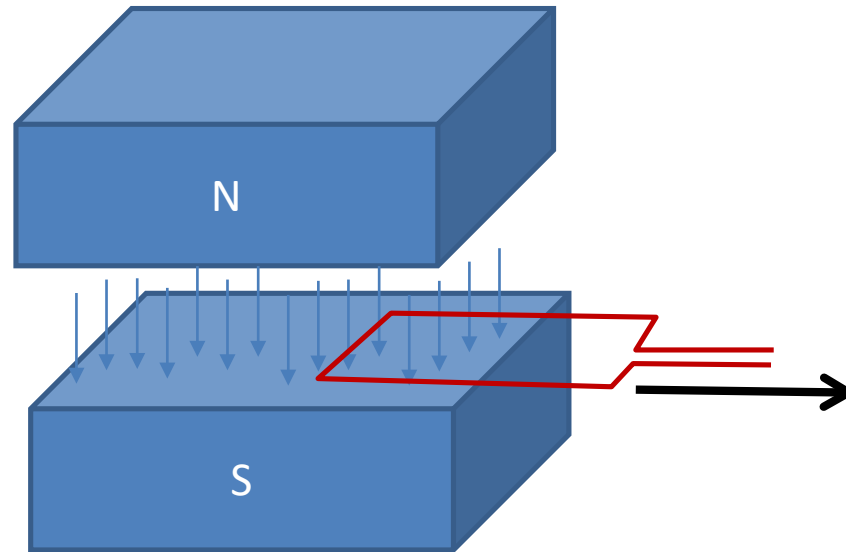
$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$$

The constants ε_0 and μ_0 are well established from the measurement of optical forces on charges and currents:

$$\mu_0 = 1.256637 \times 10^{-6} \text{N/A}^2$$

$$\varepsilon_0 = 8.854188 \times 10^{-12} \text{A}^2 \text{s}^2 / \text{N/m}^2$$

Experiments by Faraday showed that a *changing magnetic field*
Produced a force that acted in *exactly the same way as an electric field*.



This showed that the electric field and the magnetic field were *coupled together*:

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

That is: a changing magnetic field causes rotation of the electric field

These equations now look unbalanced.

$$\nabla \cdot \mathbf{E} = \frac{\rho(\mathbf{r})}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$$

In 1873, Maxwell proposed a modification of the curl equation for the magnetic field

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

What are the consequences of this?



Maxwell's equations:

$$\nabla \cdot \mathbf{E} = \frac{\rho(\mathbf{r})}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

Consider an electric and magnetic field pair in free space,
So that $\rho = 0$ and $\mathbf{J} = \mathbf{0}$ everywhere. Maxwell's equations are

$$\nabla \cdot \mathbf{E} = 0$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

Take the curl of the third equation:

The image shows two identical sets of horizontal blue lines, each set bounded by two vertical red lines. These are intended as workspace for the student to show their work for taking the curl of the third equation.

So

$$\nabla^2 \mathbf{E} = \mu_0 \varepsilon_0 \frac{\partial^2 \mathbf{E}}{\partial t^2}$$

This is a *wave equation* (we will do this in a few weeks).

It has solutions which are waves with a velocity

$$v = \frac{1}{\sqrt{\mu_0 \varepsilon_0}}$$

$$\mu_0 = 1.256637 \times 10^{-6} \text{N/A}^2$$

$$\varepsilon_0 = 8.854188 \times 10^{-12} \text{A}^2 \text{s}^2 / \text{N/m}^2$$

Plot of E-field of an electromagnetic plane wave:

